

# Transverse observables and the kinematic boundary

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Barr, BMG & Lester, arXiv:0908.3779 [hep-ph]

$m_{T(2)}$  &c

The boundary

Generalizations

Conclusions

- ▶  $m_{T(2)}$ , kinks &c.
- ▶ What does it all mean? The kinematic boundary
- ▶ Generalizations: combinatorics, non-identical decays, and the inverse of  $m_{T(2)}$

$m_{T(2)}$  &c

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In the beginning ...

# The transverse mass

$$m_0^2 = m_v^2 + m_i^2 + 2(E_v E_i - \mathbf{p}_v \cdot \mathbf{p}_i - q_v q_i)$$

- ▶  $(E, \mathbf{p}, q)$  is 4-momentum

$$m_T^2 = m_v^2 + m_i^2 + 2(e_v e_i - \mathbf{p}_v \cdot \mathbf{p}_i)$$

- ▶  $e = \sqrt{\mathbf{p} \cdot \mathbf{p} + m^2}$  is transverse energy

$$m_0 \geq m_T$$

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# Transverse Mass $m_T$

$m_{T(2)}$  &c

The boundary

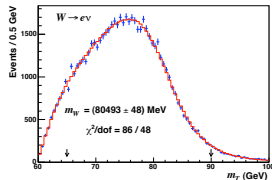
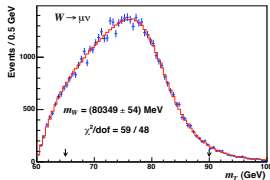
Generalizations

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$W \rightarrow l\nu$

CDF:  $m_W = 80.413 \pm 0.048$  GeV

arXiv:0708.3642



# Identical Pair Decays and $m_{T2}$

$$m_T^2 = m_V^2 + m_i^2 + 2(\mathbf{e}_V \mathbf{e}_i - \mathbf{p}_V \cdot \mathbf{p}_i)$$

▶ unobservable

$$m_{T2} = \min \max m_T$$

Lester & Summers, PLB 463 99,1999

Barr et al., J.Phys.G29:2343-2363,2003

▶ observable

$$m_{T2} \leq m_0$$

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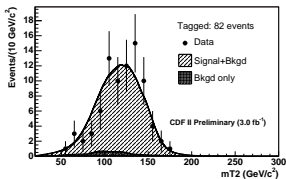
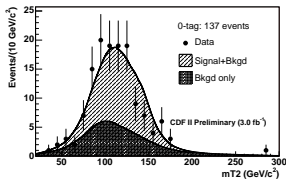
# Pair Decays and $m_{T2}$

$$\bar{t}t \rightarrow 2b2W \rightarrow 2b2l2\nu$$

Cho et al. 0804.2185

$$\text{CDF } m_{T2} \text{ only: } m_t = 167.9^{+5.6}_{-5.0} \text{ GeV}$$

CDF note 9769



# $m_{T(2)}$ and the kink

$$m_T^2 = m_V^2 + m_i^2 + 2(\mathbf{e}_V \mathbf{e}_i - \mathbf{p}_V \cdot \mathbf{p}_i)$$

- ▶  $m_T$  is unobservable if  $m_i$  unknown
- ▶ Consider  $m_T = m_T(m_i)$
- ▶ Lose boundedness but gain a kink

Choi et al., 0709.0288

BMG, JHEP 0208 051 2008

Barr, BMG & Lester, JHEP 0208 014, 2008

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$m_{T(2)}$  &c

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$m_{T(2)}$  and the kinkPair Three-body decay  $2\tilde{g} \rightarrow 2q2\bar{q}2\tilde{\chi}_1^0$ 

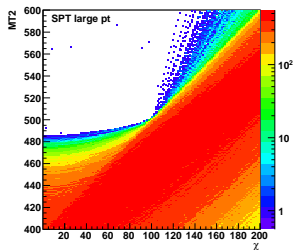
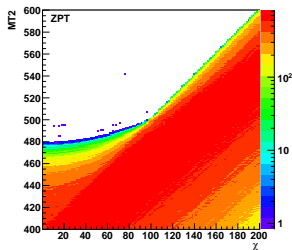
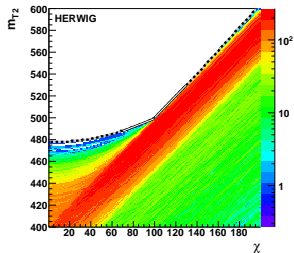
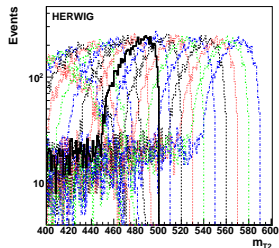
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 $m_{T(2)}$  &c

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$m_T(2)$  &c

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What does all of this mean?



# What does all of this mean?

$m_{T(2)}$  &c

The boundary

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- ▶ *ad hoc* definition of  $m_{T(2)}$
- ▶ *ad hoc* generalization to hypothesized masses
- ▶ In fact these are natural objects ...

Cheng and Han, arXiv:0810.5178

$m_{T(2)}$  &c

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# The Kinematic Boundary

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$m_{T(2)}$  is the kinematic boundary of an event

Serna, arXiv:0804.3344

Cheng and Han, arXiv:0810.5178

Kinematic constraints for an event:

- ▶  $(p_i + p_v)^2 = m_0^2$
- ▶  $p_i^2 = m_i^2$
- ▶  $\mathbf{p}_i = \mathbf{p}$

Unknowns are  $p_i$  and  $(m_i, m_0)$

$m_{T(2)}$  &c

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- ▶  $(p_i + p_v)^2 = m_0^2$
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When do these admit a solution with real momentum and real, positive energy?

Theorem: a solution exists for  $m_0 \geq m_T(m_i)$

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Theorem: a solution exists for  $m_0 \geq m_T(m_i)$

Proof: Two parts

- ▶ Part 1: Prove any solution has  $m_0 \geq m_T(m_i)$
- ▶ Part 2: Prove that there is a solution with  $m_0 = m_T(m_i)$

## Part 1

 $m_{T(2)}$  &c

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Prove any solution has  $m_0 \geq m_T(m_i)$

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- ▶  $m_T^2 \equiv m_i^2 + m_V^2 + 2(\mathbf{e}_i \mathbf{e}_V - \mathbf{p}_i \cdot \mathbf{p}_V)$
- ▶  $E_i E_V - q_i q_V \geq \mathbf{e}_i \mathbf{e}_V$ , equality at  $E_i q_V - E_V q_i = 0$
- ▶  $\implies m_0 \geq m_T(m_i)$

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# Part 2

$m_{T(2)}$  &c

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Prove  $m_0 = m_T(m_i)$  has a solution

- ▶ Find suitable  $E_i, q_i$
- ▶ Need  $\frac{q_i}{q_i^2 + e_i^2} = \frac{q_v}{E_v} \in [-1, +1]$

Proof for  $m_{T2}$  not much harder

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# The kinematic boundary

New definition of  $m_{T(2)}$  as "the kinematic boundary of an event"

So what?

- ▶ Faster algorithm for computing  $m_{T2}$

Cheng and Han, arXiv:0810.5178

- ▶  $m_{T(2)}$  is the best one can do without extra kinematics or dynamics
- ▶ Easier proofs of kinks &c.
- ▶ Generalize ...

$m_{T(2)}$  &c

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$m_{T(2)}$  &c

The boundary

**Generalizations**

Conclusions

# Generalizations

# Generalizations I: Combinatorics

$m_{T(2)}$  &c

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Conclusions

## Combinatoric ambiguities

- ▶ Branch assignments

Barr and Lester, arXiv:0708.1028

- ▶ Initial State Radiation

Alwall et al, arXiv:0905.1201

- ▶ What is the kinematic boundary?
- ▶ Union of allowed regions

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# Generalizations II: Non-identical decays

Always considered identical pair decays.

But parents/daughters need not be the same

Distinct parents:

- ▶ Squark-gluino production in MSSM

Distinct daughters

- ▶ Distinct daughters:  $> 1$  neutral particle with lifetime  $> \text{few } m$   
e.g. Neutrini, Multiple photini

Dimopoulos et al, to appear

$m_{T(2)}$  &c

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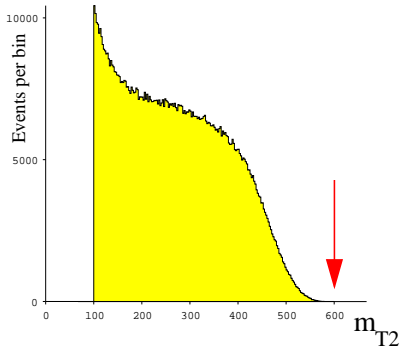
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# Generalizations II: Non-identical decays

$m_{T2}$  is not much good ...



... use the kinematic boundary?

$m_{T(2)}$  & c

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# Generalizations II: Non-identical decays

$m_{T(2)}$  &c

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## Focus on case with non-identical parents

- ▶ 3-d space of unknown masses ( $m_i, m_0, m'_0$ )
- ▶ Kinematic boundary is a surface

Extremal surface from decoupling argument ...

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Ditto for the other decay, plus

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Last equation decouples for extremal surface

$m_{T(2)}$  &c

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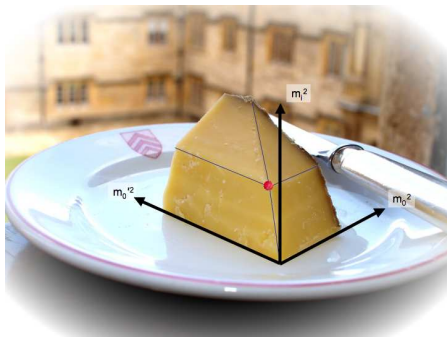
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# Generalizations II: Non-identical decays

Extremal surface



Lots of interesting kinks ...

$m_{T(2)}$  & c

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# Generalizations II: Non-identical decays

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Is there an analogue of the usual  $m_{T2}$ ?

$$m_{T2}^2(m_i, m'_i, \frac{m_0}{m'_0}) = \min \max(\frac{m'_0}{m_0} m_T^2, \frac{m_0}{m'_0} m_T'^2)$$

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## Generalizations III: The inverse of $m_{T(2)}$

$m_{T(2)}$  &c

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Use  $m_0 = m_{T(2)}(m_i)$  if you know the mass of the daughter but not the parent ...

... what about the other way round?

- ▶ Kinematic boundary the same
- ▶ Need the inverse of  $m_0 = m_{T(2)}(m_i)$



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- ▶  $m_T^2 = (\alpha_i + \alpha_v)^2$ ,  $\alpha = (\mathbf{e}, \mathbf{p})$
- ▶  $(m_T^2)^{-1} = (\alpha_0 - \alpha_v)^2$
  
- ▶  $m_{T2} = \min \max(m_T, m'_T)$
- ▶  $m_{T2}^{-1} = \max \min(m_T^{-1}, m'^{-1}_T)$

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# Conclusions

$m_{T(2)}$  & c

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- ▶  $m_{T(2)}$  & c are natural objects: define kinematic boundary
- ▶ Also the "best" objects
- ▶ Easily generalized